

Towards the minimal effective theory for leptogenesis, dark matter, and neutrino masses

Peter Maták

In collaboration with T. Blažek, J. Ramaj and M. Sabová

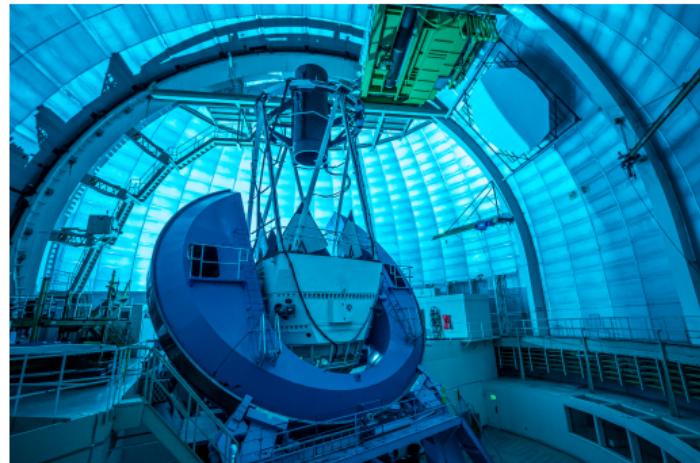
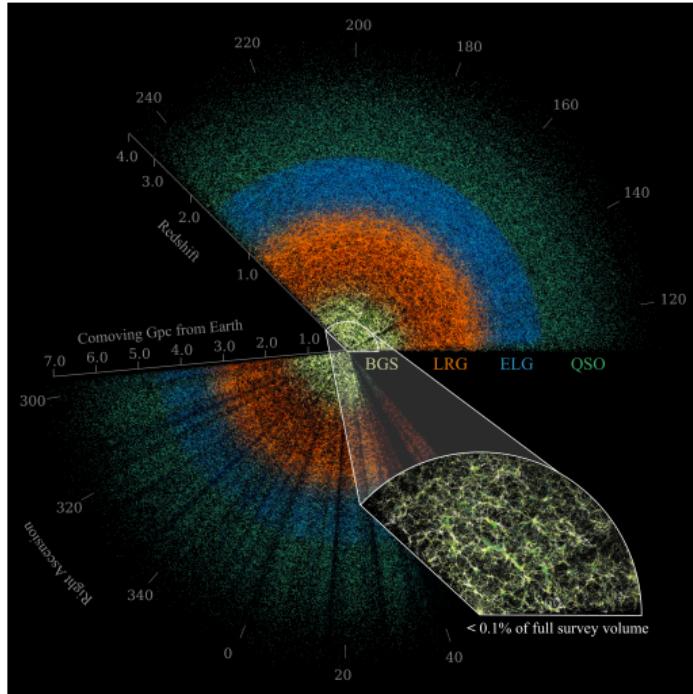


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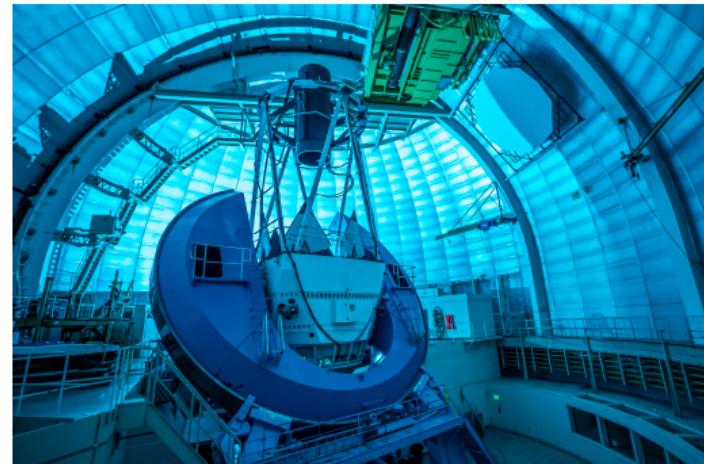
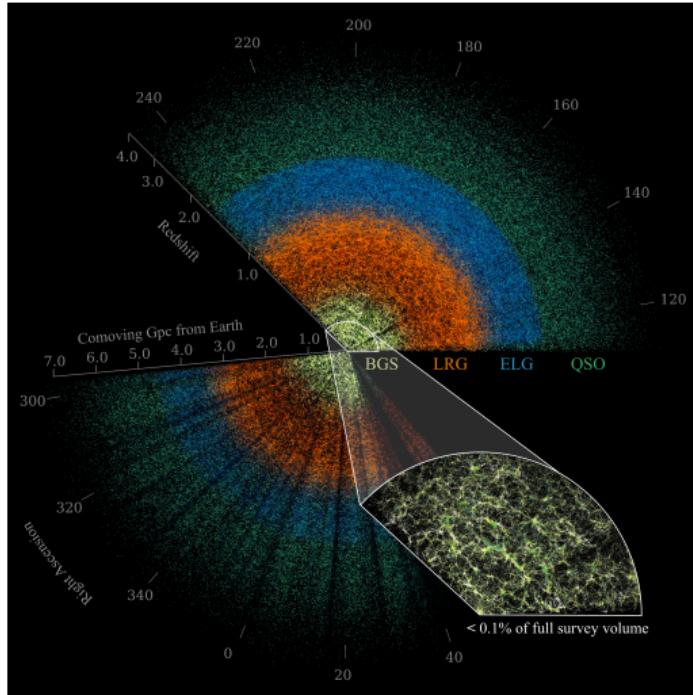
MATFYZ
CONNECTIONS

The expanding universe as we see it today



The Dark Energy Spectroscopic Instrument maps the universe by collecting spectra from millions of galaxies and quasars. Credit: Marilyn Sargent/Berkeley Lab.

The expanding universe as we see it today



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At cosmological scales, the universe is homogeneous and isotropic . . .

The expanding universe as we see it today

$$\mathbf{r}(t) = \mathbf{r}(t_0)a(t)$$

$$\left(\frac{da/dt}{a}\right)^2 = \frac{8\pi G}{3}\rho(a)$$

- ... and its only dynamics are the expansion or collapse.

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- ... and its only dynamics are the expansion or collapse.
- The universe expands, but galaxies themselves do not.
- The expansion slows massive particles down. For photons, it causes a redshift, which is not due to the Doppler effect.
- By measuring a galaxy's redshift, we gain insight into the universe's expansion rate in the past.

The expanding universe as we see it today

$$\rho(a) = \rho_{\text{crit},0} \times \left(\Omega_{\Lambda,0} + \frac{\Omega_{m,0}}{a^3} + \frac{\Omega_{r,0}}{a^4} \right)$$

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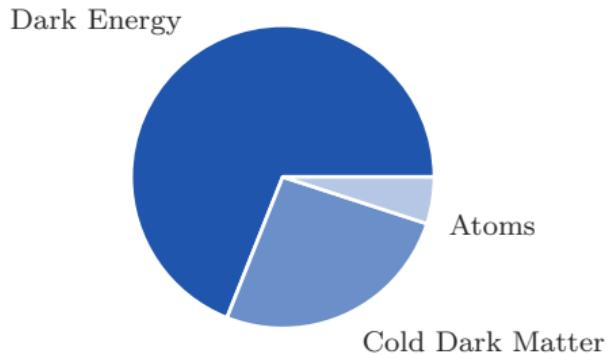
- Our universe is almost flat and has a critical density of $8.6 \times 10^{-27} \text{ kg/m}^3$.
- The expansion rate today, known as the Hubble constant, is

$$H_0 = \left. \frac{da/dt}{a} \right|_{\text{today}} = (67.7 \pm 0.4) \times \text{km s}^{-1} \text{ Mpc}^{-1}.$$

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$$\Omega_{\Lambda,0} = (68.9 \pm 0.60)\%$$

$$\left. \begin{array}{l} \Omega_{c,0} = (26.0 \pm 0.36)\% \\ \Omega_{b,0} = (4.89 \pm 0.07)\% \end{array} \right\} \Omega_{m,0} = (31.1 \pm 0.6)\%$$

$$\Omega_{r,0} = (5.38 \pm 0.15) \times 10^{-5} + \text{neutrinos}$$

S. Navas *et al.* (Particle Data Group), Phys. Rev. D 110, 030001 (2024)

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accelerates the expansion

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relative motion of galaxies,
rotational curves, ...

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mostly hydrogen

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Radiation-dominated universe

$$\rho(a) = \rho_{\text{crit},0} \times \left(\Omega_{\Lambda,0} + \frac{\Omega_{m,0}}{a^3} + \frac{\Omega_{r,0}}{a^4} \right) \approx \rho_{\text{crit},0} \times \frac{\Omega_{r,0}}{a^4} \quad \text{for } a \ll 10^{-4}$$

$$\left. \begin{array}{l} t \ll t_{\text{eq}} \approx 50000 \text{ y} \\ T \gg T_{\text{eq}} \approx 0.25 \text{ eV} \simeq 10^4 \text{ K} \end{array} \right\} a(t) \propto t^{1/2}$$

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- The universe is expanding and cooling.
- No large-scale structures have formed yet.

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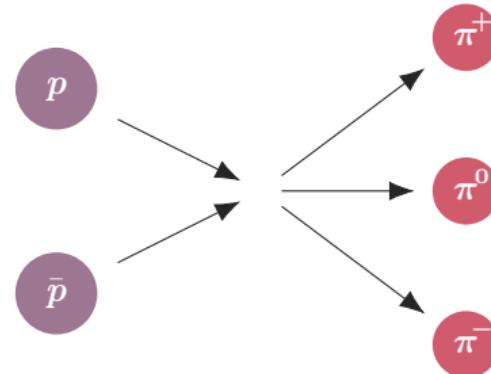
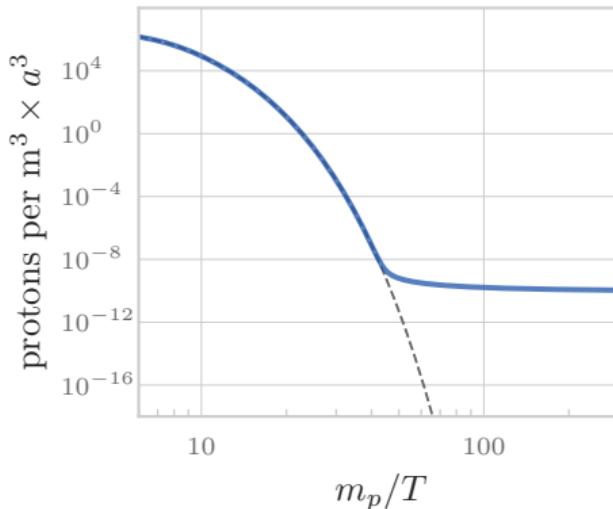
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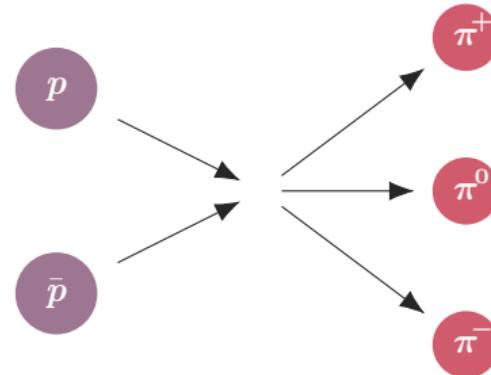
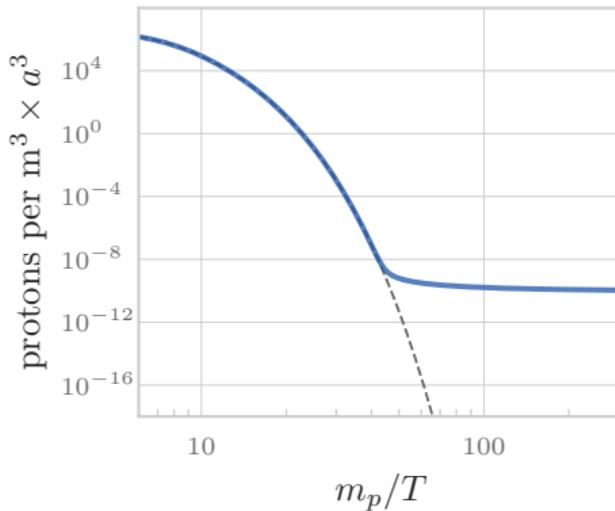
The radiation-dominated universe is a homogeneous, hot, and dense mixture of particles.

Proton anti-proton annihilation symmetric freeze-out



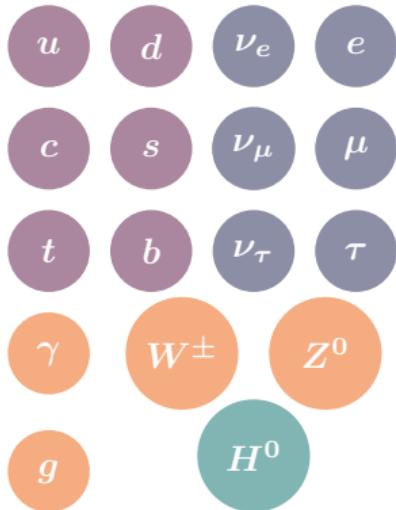
- Baryons (protons and neutrons) form at $T \approx 150$ MeV (1.7×10^{12} K)
- Annihilation and production remain in equilibrium until $T \approx 20$ MeV (2.3×10^{11} K).
- The resulting density today would be approximately 10^{-10} protons/m³.

Proton anti-proton annihilation symmetric freeze-out



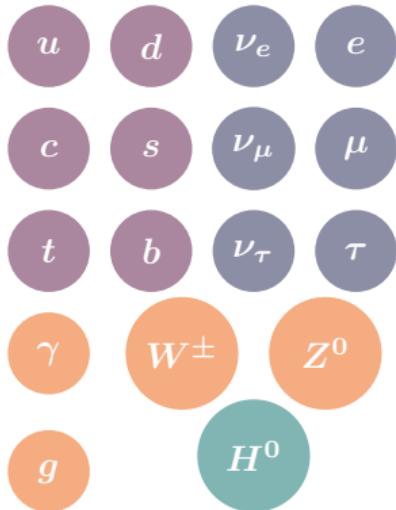
We need an excess of protons over antiprotons!
For every 10^{10} antiprotons, there must be one extra proton.

Standard Model and the early universe



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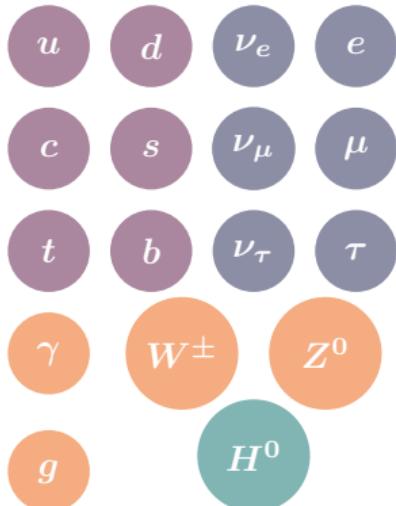
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Standard Model and the early universe

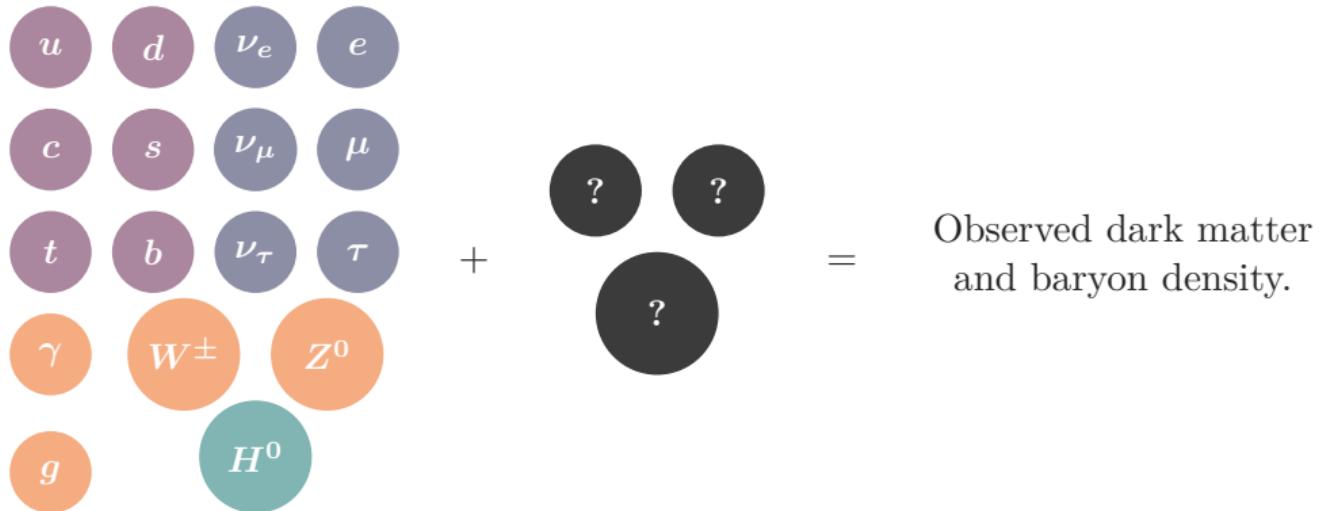


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New physics needed to explain the matter in the universe!

Standard Model and the early universe



Baryogenesis and leptogenesis

- Processes violating baryon number.
- Independent C - and CP -symmetry violation.
- Interactions out of thermal equilibrium.

A.D. Sakharov, 1991 Sov. Phys. Usp. 34 392

C : particle \leftrightarrow antiparticle, P : left \leftrightarrow right, T : past \leftrightarrow future.

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C : particle \leftrightarrow antiparticle, P : left \leftrightarrow right, T : past \leftrightarrow future.

Under reasonable assumptions (space-time symmetries, causality), any theory of particle interactions must be CPT -invariant.

J. Schwinger, Phys. Rev. 82, 914 (1951)

Baryogenesis and leptogenesis

- In **leptogenesis**, a lepton number asymmetry is generated first and then converted into a baryon asymmetry through Standard Model interactions.

M. Fukugita, T. Yanagida, Phys. Lett. B 174(1), 45 (1986)

$$\gamma_{N \rightarrow \ell H} = \frac{\# \text{ of } N \rightarrow \ell H \text{ decays}}{\text{unit volume} \times \text{unit time}}$$

- We require $\gamma_{N \rightarrow \ell H}^{\text{eq}} > \gamma_{N \rightarrow \bar{\ell} \bar{H}}^{\text{eq}}$ and this necessarily implies $\gamma_{\ell H \rightarrow \bar{\ell} \bar{H}}^{\text{eq}} > \gamma_{\bar{\ell} \bar{H} \rightarrow \ell H}^{\text{eq}}$.
- In thermal equilibrium, decay and scattering asymmetries cancel each other, so a departure from equilibrium is required.
- Stable dark matter candidate still missing.

Leptogenesis and dark matter in a minimal model

Top-down approach

- Well-motivated (neutrino masses, naturalness).
- Built upon new symmetries.
- Examples: SUSY, GUTs, extended Higgs sectors, ...

Bottom-up approach

- Minimal in terms of new particles and interactions.
- Effective theory below some energy scale Λ (power series in T/Λ).

How minimal can an effective theory be?

- Two new particles (a heavy neutral fermion f , a light neutral scalar S) stabilized by a \mathbb{Z}_2 symmetry.
- One portal operator and the Weinberg operator (leads to neutrino masses),

$$L_{\text{eff}} = \frac{\lambda}{\Lambda} S \bar{f} P_L \ell H + \frac{\lambda'}{\Lambda} H \bar{\ell}^c P_L \ell H + \text{H.c.}$$

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T. Blažek, P. Maták, J. Ramaj, M. Sabová, Eur. Phys. J. C 85 (2025) 801

Thank you for your attention!